

# Neutrino Oscillation

(1)

## A) Vacuum Oscillation

$$\left. \begin{aligned} \nu^{(f)} &= \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} \rightarrow \text{flavor state} \\ \nu^{(p)} &= \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} \rightarrow \text{mass state} \end{aligned} \right\} \text{orthonormal}$$

$$\nu^{(f)} = U \nu^{(p)}$$

$$U = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \rightarrow \text{mixing matrix (Unitary)}$$

### Evolution equation

$$i \frac{d}{dt} \nu^{(p)}(x) = H \nu^{(p)}(x) \quad \left| \quad H = \begin{pmatrix} E_1 & 0 \\ 0 & E_2 \end{pmatrix} \rightarrow \text{diagonalized} \right.$$

$$\Rightarrow i \frac{d}{dx} \nu^{(p)}(x) = H \nu^{(p)}(x) \quad \left. \begin{array}{l} \text{Replacing } t \text{ by } x \\ c = \hbar = 1 \end{array} \right|$$

In the solar neutrino problem energy  $\rightarrow$  MeV range  
Lab. exp.  $\Rightarrow m_{\nu_e} < 9.3 \text{ eV} \quad m_{\nu_\mu} < 250 \text{ KeV}$

$$\Rightarrow m_\alpha \ll E_\alpha \quad (\alpha = 1, 2)$$

$$\Rightarrow E_\alpha \equiv \sqrt{|\vec{p}|^2 + m_\alpha^2} \simeq |\vec{p}| + \frac{m_\alpha^2}{2|\vec{p}|}$$

$$H = |\vec{p}| + \frac{1}{2|\vec{p}|} \begin{pmatrix} m_1^2 & 0 \\ 0 & m_2^2 \end{pmatrix} \\ = \left( |\vec{p}| + \frac{m_1^2 + m_2^2}{4|\vec{p}|} \right) - \frac{\Delta}{4|\vec{p}|} \sigma_3$$

$$\left. \begin{aligned} \Delta &\equiv m_2^2 - m_1^2 \\ \sigma_1 &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ \sigma_2 &= \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \\ \sigma_3 &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{aligned} \right|$$

$$i \frac{d}{dx} \psi(x) = H \psi(x)$$

$$\Rightarrow i \frac{d}{dx} [U^\dagger \psi(x)] = H U^\dagger \psi(x)$$

$$\Rightarrow i \frac{d}{dx} \psi(x) = U H U^\dagger \psi(x) = H' \psi(x)$$

$$H' = U H U^\dagger = - \begin{pmatrix} -c_{2\theta} & \sin 2\theta \\ \sin 2\theta & c_{2\theta} \end{pmatrix}$$

$$\Rightarrow H' = |\bar{F}| + \frac{m_1^2 + m_2^2}{4|\bar{F}|} + \frac{\Delta}{4|\bar{F}|} U \sigma_3 U^\dagger$$

$$\Rightarrow H' = |\bar{F}| + \frac{m_1^2 + m_2^2}{4|\bar{F}|} + \frac{\Delta}{4|\bar{F}|} \begin{pmatrix} -c_{2\theta} & \sin 2\theta \\ \sin 2\theta & c_{2\theta} \end{pmatrix}$$

$$\Rightarrow \begin{pmatrix} H'_{11} & H'_{12} \\ H'_{12} & H'_{22} \end{pmatrix} = \begin{pmatrix} S - \frac{\Delta}{4|\bar{F}|} c_{2\theta} & \frac{\Delta}{4|\bar{F}|} \sin 2\theta \\ \frac{\Delta}{4|\bar{F}|} \sin 2\theta & S + \frac{\Delta}{4|\bar{F}|} c_{2\theta} \end{pmatrix}$$

$$S = |\bar{F}| + \frac{m_1^2 + m_2^2}{4|\bar{F}|} \quad \Delta = m_2^2 - m_1^2$$

$$\Rightarrow H'_{22} - H'_{11} = \frac{\Delta}{2|\bar{F}|} c_{2\theta} \quad H'_{12} = \frac{\Delta}{4|\bar{F}|} \sin 2\theta$$

$$\boxed{\tan 2\theta = \frac{2H'_{12}}{H'_{22} - H'_{11}}}$$

The term in  $H'$ , which is proportional to unit matrix, gives an overall phase to the solution. Such a term does not affect the mixing angle as is seen above.

Thus the term  $S$  is dropped

$$i \frac{d}{dx} \psi^{(s)} = H' \psi^{(s)}$$

Solution  $\rightarrow \psi^{(s)}(x) = \exp(-iH'x) \psi^{(s)}(0)$

$$H' = \frac{\Delta}{4E} [O_1 \sin 2\theta - O_3 \cos 2\theta]$$

$$\Rightarrow \psi^{(s)}(x) = \exp\left[-\frac{i\Delta}{4E} x (O_1 \sin 2\theta - O_3 \cos 2\theta)\right] \psi^{(s)}(0)$$

$$\Rightarrow \psi^{(s)}(x) = \left[ \cos \frac{\Delta}{4E} x - i (O_1 \sin 2\theta - O_3 \cos 2\theta) \frac{\sin \frac{\Delta}{4E} x}{\frac{\Delta}{4E}} \right] \psi^{(s)}(0)$$

$$\Rightarrow P_{\nu_e \nu_e}(x) = |\langle \nu_e | \psi^{(s)}(x) \rangle|^2 = \sin^2 2\theta \sin^2\left(\frac{\Delta}{4E} x\right)$$

$$\Rightarrow P_{\nu_e \nu_e}(x) = 1 - P_{\nu_e \nu_\mu}(x)$$

$\Rightarrow$  The probability of finding  $\nu_e$  is less than unity

$\Rightarrow$  The flux depletion observed in the solar neutrino experiment.